# Weak Asymptotics for the Generating Polynomials of the Stirling Numbers of the Second Kind 

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For the horizontal generating functions $P_{n}(z)=\sum_{k=1}^{n} S(n, k) z^{k}$ of the Stirling numbers of the second kind, a weak asymptotic is established, as $n \rightarrow \infty$. The distribution function $F_{n}$ of the zeros of $Q_{n}(z)=P_{n}(n z)$ is investigated and by using the Stieltjes transformation, the limit of $F_{n}$ in the sense of weak convergence is deduced. © 2001 Academic Press
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## 1. INTRODUCTION AND SUMMARY

In this paper we deduce the asymptotic distribution function of the zeros for the horizontal generating function of the Stirling numbers of the second kind $S(n, k)$, which are defined by the following double generating function (see [2, p. 50]):

$$
\begin{equation*}
\exp \left\{z\left(e^{u}-1\right)\right\}=: 1+\sum_{1 \leqslant k \leqslant n<\infty} S(n, k) \frac{u^{n}}{n!} z^{k}, \quad z, u \in \mathbb{C} . \tag{1.1}
\end{equation*}
$$

The horizontal generating functions $P_{n}(z)$ are the coefficients of the following power series:

$$
\exp \left\{z\left(e^{u}-1\right)\right\}=: 1+\sum_{n=1}^{\infty} \frac{P_{n}(z)}{n!} u^{n}, \quad z, u \in \mathbb{C} .
$$

Concerning the zeros of $P_{n}$ there is one result: The zeros are simple, real, and not greater than 0 (see [2, p. 271]). In [5] two asymptotic expansions for $Q_{n}(z):=P_{n}(n z)$ are deduced by means of the saddle point method. With $\Phi: \mathbb{C} \backslash\{-e, 0\} \rightarrow \mathscr{A} \cup \Gamma_{+}$, which maps $\mathbb{C} \backslash[-e, 0]$ conformally on 218
$\mathscr{A}:=\{w \in \mathbb{C} \backslash\{0\}: w>-1$ or $w=a+i b, a>-b \cot b, b \in(-\pi, \pi) \backslash\{0\}\}$, and $(-e, 0)$ one-one on $\Gamma_{+}:=\{w \in \partial \mathscr{A}: \mathfrak{J}(w)>0\}$, the following results are obtained:
(i) With $\phi \in(0, \pi)$ there is the oscillating asymptotics

$$
Q_{n}(x(\phi))=k_{n}(\phi)\left(\sin \left(n\left(\pi-\phi+\frac{\sin ^{2} \phi}{\phi}\right)+\eta(\phi)\right)+\mathcal{O}\left(\frac{1}{n}\right)\right)
$$

with $k_{n}(\phi)>0, \eta(\phi)$ bounded by $\frac{\pi}{2}$ and $\pi, x(\phi) \in(-e, 0), x(\phi) w(\phi) e^{w(\phi)}$ $=1$,

$$
\begin{equation*}
x(\phi)=-\frac{\sin \phi}{\phi} e^{\phi \cot \phi} \quad \text { and } \quad \Phi(x(\phi))=w(\phi):=\frac{\phi}{\sin \phi} e^{i(\pi-\phi)} . \tag{1.2}
\end{equation*}
$$

(ii) With $z \in \mathbb{C} \backslash[-e, 0], w=\Phi(z)$ and $z w e^{w}=1$, it holds that

$$
Q_{n}(z)=\frac{n!}{\sqrt{2 \pi n}} \frac{1}{w^{n}} \exp \left\{\frac{n}{w}\left(1-e^{-w}\right)\right\}(1+w)^{-1 / 2}\left(1+\mathcal{O}\left(\frac{1}{n}\right)\right),
$$

where the $\mathcal{O}$-term holds uniformly on every compact subset of $\mathbb{C} \backslash[-e, 0]$. Examining these results, we can presume that nearly all zeros of $Q_{n}$ are in $[-e, 0]$, because for sufficiently large $n$ the asymptotic in (ii) is not equal to 0 . Therefore, we discuss the position of the zeros by investigating the distribution function of the zeros for $Q_{n}(z)$ :

$$
F_{n}(\xi):=\frac{1}{n} N_{n}(\xi), \quad N_{n}(\xi):=\left|\left\{x \leqslant \xi: Q_{n}(x)=0\right\}\right|, \quad n \in \mathbb{N} .
$$

By using the Stieltjes transformation, we get a weak asymptotic; i.e., we prove the weak convergence of the sequence $F_{n}$ to a distribution function $F$ of the following type,

$$
F(x)= \begin{cases}0, & x \leqslant-e \\ 1+\frac{1}{\pi}\left(\mathfrak{J}\left(\frac{1}{w}\right)-\operatorname{ph}(w)\right), & x \in(-e, 0) \\ 1, & x \geqslant 0,\end{cases}
$$

with $\operatorname{ph}(w) \in(0, \pi), w \in \partial \mathscr{A}, x w e^{w}=1$. With the above parametrization (1.2), it holds further that

$$
F\left(-\frac{\sin \phi}{\phi} e^{\phi \cot \phi}\right)=\frac{1}{\pi}\left(\phi-\frac{\sin ^{2} \phi}{\phi}\right), \quad \phi \in(0, \pi) .
$$

## 2. WEAK ASYMPTOTICS

To investigate the position of the zeros of $Q_{n}$, which are all simple, real, and not greater than 0 (see [2, p.271]), we define

$$
\begin{aligned}
& Q_{n}(z)=n^{n} \prod_{v=1}^{n}\left(z-x_{n, v}\right) \\
& \quad \text { with }-\infty<x_{n, 1}<\cdots<x_{n, n}=0, \quad n \in \mathbb{N}, \\
& N_{n}(\xi):=\left|\left\{v \in\{1, \ldots, n\}: x_{n, v} \leqslant \xi\right\}\right|, \quad n \in \mathbb{N}, \quad \text { and } \\
& F_{n}(\xi):=\frac{1}{n} N_{n}(\xi), \quad n \in \mathbb{N} .
\end{aligned}
$$

For computing the limit probability distribution $F$ in the sense of weak convergence, we observe the logarithmic derivative

$$
\begin{align*}
h_{F_{n}}(z) & :=\frac{1}{n} \frac{Q_{n}^{\prime}(z)}{Q_{n}(z)}=\frac{1}{n} \sum_{v=1}^{n} \frac{1}{z-x_{n, v}} \\
& =\int_{-\infty}^{0} \frac{1}{z-t} d F_{n}(t), \quad z \in \mathbb{C} \backslash(-\infty, 0],
\end{align*}
$$

which represents the Stieltjes transform of $F_{n}$. The following theorem supplies the limit of this sequence.

Theorem 2.1. With the above mentioned notations, $h_{F_{n}}$ and $F_{n}$ hold:
(i) $\quad h_{F_{n}}(z)$ converges to $h(z):=e^{\Phi(z)}-1$ (compactly) on $\mathbb{C} \backslash[-e, 0]$.
(ii) The sequence of distributions $F_{n}$ converges weakly to a distribution $F$; i.e., if $F$ is continuous at $\xi \in \mathbb{R}$, then $\lim _{n \rightarrow \infty} F_{n}(\xi)=F(\xi)$ and:

$$
\int_{-\infty}^{0} \frac{1}{z-t} d F(t)=\int_{-e}^{0} \frac{1}{z-t} d F(t)=h(z), \quad z \in \mathbb{C} \backslash[-e, 0] .
$$

Proof. (i) Let $K \subset \mathbb{C} \backslash[-e, 0]$ be compact, $z \in K$, and

$$
G_{n}(z):=\frac{n!}{\sqrt{2 \pi n}} \frac{1}{\Phi(z)^{n}} \exp \left\{n\left(\frac{1}{\Phi(z)}-z\right)\right\}(1+\Phi(z))^{-1 / 2} .
$$

Because $Q_{n} / G_{n}$ converges uniformly to 1 on $K$ (see [5, Theorem 3.3]), it follows from (2.1) that

$$
\begin{aligned}
\lim _{n \rightarrow \infty} h_{F_{n}}(z)= & \lim _{n \rightarrow \infty} \frac{1}{n} \frac{G_{n}^{\prime}(z)}{G_{n}(z)} \\
= & \lim _{n \rightarrow \infty} \frac{1}{n}\left(-n \frac{\Phi^{\prime}(z)}{\Phi(z)}-n\left(\frac{\Phi^{\prime}(z)}{\Phi(z)^{2}}+1\right)\right. \\
& \left.-\frac{1}{2}(1+\Phi(z))^{-1} \Phi^{\prime}(z)\right) \\
= & -\frac{\Phi^{\prime}(z)}{\Phi(z)}-\frac{\Phi^{\prime}(z)}{\Phi(z)^{2}}-1 .
\end{aligned}
$$

The uniform convergence follows from the compact convergence in [5, Theorem 3.3] and because $\Phi$ and $\Phi^{\prime}$ are bounded on $K$. Let $\Psi$ be the inverse function of $\Phi, \Psi(w)=\left(w e^{w}\right)^{-1}, w \in \mathscr{A}$ (see [5, Lemma 2.3]): then we obtain that $\Phi^{\prime}(z)=\left(\Psi^{\prime}(\Phi(z))\right)^{-1}, \Psi^{\prime}(w)=e^{-w}\left(-w^{-2}-w^{-1}\right)$, and further that

$$
\begin{aligned}
\lim _{n \rightarrow \infty} h_{F_{n}}(z) & =\frac{e^{\phi(z)}}{-\Phi(z)^{-2}-\Phi(z)^{-1}}\left(-\frac{1}{\Phi(z)}-\frac{1}{\Phi(z)^{2}}\right)-1 \\
& =\frac{e^{\phi(z)}}{1+\Phi(z)}(\Phi(z)+1)-1=e^{\phi(z)}-1=h(z) .
\end{aligned}
$$

(ii) By using the theorem of Grommer and Hamburger (see [10, p. 175]) and Lemma 1 in [6, p. 383], the proof is completed if $\mathfrak{R}(i y h(i y)) \rightarrow 1$ for $0<y \rightarrow \infty$. This follows from [5], Lemma 2.3:

$$
\lim _{z \rightarrow \infty} z h(z)=\lim _{w \rightarrow 0} \Psi(w) h(\Psi(w))=\lim _{w \rightarrow 0} \frac{1}{w} e^{-w}\left(e^{w}-1\right)=1 .
$$

Since $h$ is analytic on $\mathbb{C} \backslash[-e, 0]$, the lower boundary in the integral representation $-\infty$ can be replaced by $-e$.

Now we focus on determining the distribution $F$. To use some wellknown results concerning the inversion of Stieltjes transforms, we make the following assumption:

$$
\begin{equation*}
F \in C_{1}(-e, 0), F^{\prime}(t)=f(t), t \in(-e, 0) . \tag{2.2}
\end{equation*}
$$

Thereby, we can apply an inversion formula for the Stieltjes transform

$$
\begin{equation*}
h(z)=e^{\Phi(z)}-1=\int_{-e}^{0} \frac{f(t)}{z-t} d t, \quad z \in \mathbb{C} \backslash[-e, 0] . \tag{2.3}
\end{equation*}
$$

Theorem 2.2 The function $f$ with the Stieltjes transform $h$ and the distribution $F$ hold: ${ }^{1}$
(i) $f(t)=\frac{1}{\pi} \mathfrak{J}(\exp (\Phi(t))), \quad t \in(-e, 0)$,
(ii) $F(x)= \begin{cases}0, & x \leqslant-e \\ 1+\frac{1}{\pi}\left(\mathfrak{J}\left(\frac{1}{\Phi(x)}\right)-\operatorname{ph}(\Phi(x))\right), & x \in(-e, 0), \operatorname{ph}(\Phi(x)) \in(0, \pi) \\ 1, & x \geqslant 0,\end{cases}$
(iii) $F(x(\phi))=\frac{1}{\pi}\left(\phi-\frac{\sin ^{2} \phi}{\phi}\right)$,

$$
x(\phi)=-\frac{\sin \phi}{\phi} e^{\phi \cot \phi}, \quad \phi \in(0, \pi) .
$$

Proof. (i) Let $t$ be in $(-e, 0)$; then by using Theorem 14.1 in [11, p. 126]:

$$
\begin{equation*}
f(t)=\lim _{y \rightarrow 0_{+}} \frac{h(t-i y)-h(t+i y)}{2 \pi i} . \tag{2.4}
\end{equation*}
$$

To compute $h(t-i 0)$ let $z \notin \mathbb{R}, \mathfrak{R}(z) \in(-e, 0), z w e^{w}=1, w=x+i y, y \neq 0$, and $w \in \mathscr{A}$; then we have:

$$
\begin{aligned}
\mathfrak{J}(z) & =\mathfrak{J}\left(\frac{1}{x+i y} e^{-x-i y}\right)=\mathfrak{J}\left(\frac{x-i y}{x^{2}+y^{2}} e^{-x} e^{-i y}\right) \\
& =\frac{e^{-x}}{x^{2}+y^{2}}(-x \sin y-y \cos y) \\
& =\frac{e^{-x}}{x^{2}+y^{2}} \sin y(-x-y \cot y) .
\end{aligned}
$$

So by the definition of $\mathscr{A}$ we can conclude that $\mathfrak{J}(z)$ is less than or greater than 0 if and only if $y$ is greater than or less than 0 , respectively. Thus (2.4) becomes:

$$
f(t)=\frac{1}{2 \pi i}(\exp (\Phi(t))-\exp (\overline{\Phi(t)}))=\frac{1}{\pi} \mathfrak{J}(\exp (\Phi(t)))
$$

[^0]Because $f$ has been determined by the help of (2.2) we still have to verify that $f$ complies with (2.3). That can be easily considered by the substitutions $t=\Psi(w), w \in \Gamma_{+}$, and $t=\Psi(w), w \in \Gamma_{-}$, and the use of the residue theorem.
(ii) By Theorem 2.1 and (2.2) there is a $c$ in $[0,1)$ satisfying:

$$
F(x)= \begin{cases}0, & x<-e  \tag{2.5}\\ c, & x=-e \\ c+\int_{-e}^{x} f(t) d t, & x \in(-e, 0) \\ 1, & x \geqslant 0 .\end{cases}
$$

With $x$ in $(-e, 0)$ and $t=\Psi(w)=\left(w e^{w}\right)^{-1}$, it follows that

$$
\begin{aligned}
\int_{-e}^{x} \frac{1}{\pi} \mathfrak{J}(\exp (\Phi(t))) d t & =\frac{1}{\pi} \mathfrak{J}\left(\int_{-1, w \in \Gamma_{+}}^{\Phi(x)} e^{w} e^{-w}\left(-\frac{1}{w^{2}}-\frac{1}{w}\right) d w\right) \\
& =\frac{1}{\pi} \mathfrak{J}\left(\frac{1}{\Phi(x)}+1-\ln (\Phi(x))+\ln (-1)\right) \\
& =1+\frac{1}{\pi}\left(\mathfrak{J}\left(\frac{1}{\Phi(x)}\right)-\operatorname{ph}(\Phi(x))\right),
\end{aligned}
$$

which has the limit 0 for $x \rightarrow-e_{+}$and 1 for $x \rightarrow 0_{-}$. Hence, $c$ must be equal to 0 and the proof is completed.
(iii) Let $\Phi$ be in $(0, \pi)$; then due to (ii) and (1.2) we may write:

$$
\begin{aligned}
F(x(\phi)) & =1+\frac{1}{\pi}\left(\mathfrak{J}\left(\frac{1}{\Phi(x(\phi))}\right)-\operatorname{ph}(\Phi(x(\phi)))\right) \\
& =1+\frac{1}{\pi}\left(\mathfrak{J}\left(\frac{\sin \phi}{\phi} e^{-i(\pi-\phi)}\right)-\operatorname{ph}\left(\frac{\phi}{\sin \phi} e^{i(\pi-\phi)}\right)\right) \\
& =1-\frac{1}{\pi} \frac{\sin ^{2} \phi}{\phi}-1+\frac{\phi}{\pi}=\frac{1}{\pi}\left(\phi-\frac{\sin ^{2} \phi}{\phi}\right) .
\end{aligned}
$$

Besides the access to the density $f$ with the help of the Stieltjes transform, there is also the possibility of using an inversion formula for the Mellin transform $g(s)$ of $f_{*}(t):=f(-t)$ :

$$
\begin{equation*}
g(s):=\int_{0}^{e} f_{*}(t) t^{s-1} d t=\int_{-e}^{0} f(t)(-t)^{s-1} d t, \quad \mathfrak{R}(s)>1 . \tag{2.6}
\end{equation*}
$$

Indeed, that does not lead to the above mentioned form of $f$, but we can obtain some interesting results by using this method.

Theorem 2.3. Let $x$ be in $(-e, 0), \mathfrak{R}(s)>1$ and $\ln (s-1)=\ln |s-1|+$ $i \operatorname{ph}(s-1), \operatorname{ph}(s-1) \in(-\pi, \pi)$, then:

$$
\begin{equation*}
g(s)=\frac{(s-1)^{(s-1)}}{\Gamma(s+1)} \tag{i}
\end{equation*}
$$

(ii) $f(x)=\frac{1}{2 \pi i} \int_{c-i \infty}^{c+i \infty} \frac{(s-1)^{s-1}}{\Gamma(s+1)}(-x)^{-s} d s, \quad c>1$,
(iii) $f(x)=\frac{1}{-x \pi^{2}} \int_{0}^{\infty} \frac{(-x)^{r}}{r^{r}} \sin ^{2}(\pi r) \frac{\Gamma(r)}{1-r} d r$.

Proof. (i) To compute $g(s)$ we observe that for $|z|>e$ :

$$
\begin{aligned}
e^{\Phi(z)}-1 & =\int_{-e}^{0} \frac{f(t)}{z-t} d t=\frac{1}{z} \int_{-e}^{0} f(t) \sum_{k=0}^{\infty}\left(\frac{t}{z}\right)^{k} d t \\
& =\sum_{k=0}^{\infty} \frac{(-1)^{k}}{z^{k+1}} \int_{0}^{e} f_{*}(t) t^{k} d t .
\end{aligned}
$$

On the other hand we can develop $e^{\phi(z)}-1$ into a Bürmann Lagrange series (see [8, pp. 124-125]):

$$
e^{\Phi(z)}-1=\sum_{k=1}^{\infty} \frac{1}{z^{k}} \frac{1}{k!}\left[\frac{d^{k-1} e^{x}\left(e^{-x}\right)^{k}}{d x^{k-1}}\right]_{x=0}=\sum_{k=0}^{\infty} \frac{(-1)^{k} k^{k}}{(k+1)!} \frac{1}{z^{k+1}} .
$$

By the identity theorem for Laurent series, an identity theorem for analytic functions (Theorem 5.81 in [9, p. 186]), and by analytic continuation the claimed form of $g(s)$ is proved.
(ii) By using an inversion formula for $g(s)$ (see [4, p. 409]) it follows that

$$
\begin{equation*}
f(x)=f_{*}(-x)=\frac{1}{2 \pi i} \int_{c-i \infty}^{c+i \infty} \frac{(s-1)^{s-1}}{\Gamma(s+1)}(-x)^{-s} d s, \quad c>1 . \tag{2.7}
\end{equation*}
$$

(iii) To prove (iii) we look at a modified path of integration (Fig. 1) with $\tilde{c}=c-1, d>0, R:=|\tilde{c}-i d|, \alpha_{-}:=\operatorname{ph}(\tilde{c}-i d) \in\left(-\frac{\pi}{2}, 0\right), \alpha_{+}=\operatorname{ph}(\tilde{c}+i d)$ $=-\alpha_{-}$, and:

$$
\begin{aligned}
\gamma_{1} & :=\{s \in \mathbb{C}: s=\tilde{c}+i y, y \in[-d, d]\}, \\
\gamma_{2} & :=\left\{s \in \mathbb{C}: s=R e^{-i \psi}, \psi \in\left[\alpha_{+}, \pi\right]\right\}, \\
\gamma_{3} & :=\left\{s \in \mathbb{C}: s=r e^{-i \pi}, s \in[-R,-1 / R]\right\}, \\
\gamma_{4} & :=\left\{s \in \mathbb{C}: s=R^{-1} e^{i \psi}, \psi \in[-\pi, \pi]\right\}, \\
\gamma_{5} & :=\left\{s \in \mathbb{C}: s=r e^{i \pi}, r \in[1 / R, R]\right\}, \\
\gamma_{6} & :=\left\{s \in \mathbb{C}: s=\operatorname{Re} e^{i(\pi-\psi)}, \psi \in\left[0, \pi-\alpha_{+}\right]\right\} .
\end{aligned}
$$



FIG. 1. The path of integration for Theorem 2.3(iii).
By Cauchy's theorem and analytic continuation of $g$ it follows that

$$
\begin{align*}
\int_{c-i \infty}^{c+i \infty} & g(s)(-x)^{-s} d s \\
\quad= & \lim _{d \rightarrow \infty} \int_{\gamma_{2}+\gamma_{3}+\gamma_{4}+\gamma_{5}+\gamma_{6}} g(s+1)(-x)^{-s-1} d s \tag{2.8}
\end{align*}
$$

For investigating $\int_{\gamma_{2}}$ let $s=R e^{-i \psi}, R>1, \psi \in\left[\alpha_{+}, \pi\right]$; then by the functional equation of the Gamma function (see [1, p. 48]) and Stirling's formula (see [7, p. 294]) it holds that

$$
\begin{aligned}
g(s+ & 1)(-x)^{-s-1} \\
= & \frac{s^{s}}{\Gamma(s+2)}(-x)^{-s-1} \\
= & \frac{s^{s}}{s+1} \frac{\sin (-\pi s)}{\pi} \Gamma(-s)(-x)^{-s-1} \\
= & \frac{s^{s}}{s+1} \frac{\sin (-\pi s)}{\pi} e^{s}(-s)^{-s} \sqrt{\frac{2 \pi}{-s}}(-x)^{-s-1}\left(1+\mathcal{O}\left(\frac{1}{-s}\right)\right) \\
= & \frac{e^{s} \sqrt{2}}{(-x)^{s+1}(s+1) \sqrt{R \pi}} \frac{e^{-i \pi s}-e^{i \pi s}}{2 i} \\
& \times e^{-(1 / 2) i(-\psi+\pi)} e^{s \ln (R)-i s \psi} e^{-s \ln (R)-i s(-\psi+\pi)}\left(1+\mathcal{O}\left(\frac{1}{-s}\right)\right) \\
= & \left(\frac{e}{-x}\right)^{s} \frac{-1}{-x(s+1) \sqrt{2 \pi s}}\left(e^{-2 i \pi s}-1\right)\left(1+\mathcal{O}\left(\frac{1}{-s}\right)\right)
\end{aligned}
$$

Since $\mathfrak{J}(s) \leqslant 0$ and $\mathfrak{R}(s) \leqslant \tilde{c}, \int_{\gamma_{2}}$ is satisfying:

$$
\begin{align*}
& \left|\int_{\gamma_{2}} g(s+1)(-x)^{-s-1} d s\right| \\
& \quad \leqslant\left(\frac{e}{-x}\right)^{\tilde{c}} \frac{1}{|x|(R-1) \sqrt{2 \pi R}} 2 \pi R\left(1+\mathcal{O}\left(\frac{1}{R}\right)\right) \rightarrow 0, \quad d \rightarrow \infty . \tag{2.9}
\end{align*}
$$

By analogous argumentation it follows that

$$
\begin{equation*}
\left|\int_{\gamma_{6}} g(s+1)(-x)^{-s-1} d s\right| \rightarrow 0, \quad d \rightarrow \infty \tag{2.10}
\end{equation*}
$$

To estimate $\int_{\gamma_{4}}$ let $s=R^{-1} e^{i \psi}, \psi \in[-\pi, \pi]$. Since $(\Gamma(s+2))^{-1} \rightarrow 1$ and $\left|s^{s}\right| \rightarrow 1$ as $d \rightarrow \infty$, it follows that

$$
\begin{equation*}
\left|\int_{\gamma_{4}} g(s+1)(-x)^{-s-1} d s\right| \leqslant \max _{s \in \gamma_{4}}\left|\frac{s^{s}(-x)^{-s-1}}{\Gamma(s+2)}\right| \frac{2 \pi}{R} \rightarrow 0 \tag{2.11}
\end{equation*}
$$

as $d \rightarrow \infty$. Altogether, we conclude from (2.7), (2.8), (2.9), (2.10), and (2.11) with the substitutions $s=r e^{-i \pi}$ and $s=r e^{i \pi}$ :

$$
\begin{aligned}
f(x)= & \frac{1}{2 \pi i}\left(\int_{\infty}^{0} g\left(r e^{-i \pi}+1\right)(-x)^{-r e^{-i \pi}-1} e^{-i \pi} d r\right. \\
& \left.+\int_{0}^{\infty} g\left(r e^{i \pi}+1\right)(-x)^{-r e^{i \pi}-1} e^{i \pi} d r\right) \\
= & \frac{1}{-2 \pi i x} \int_{0}^{\infty} \frac{(-x)^{r}}{\Gamma(2-r)}\left(\left(r e^{-i \pi}\right)^{r e^{-i \pi}}-\left(r e^{i \pi}\right)^{r e i \pi}\right) d r \\
= & \frac{1}{-2 \pi i x} \int_{0}^{\infty} \frac{(-x)^{r}}{\Gamma(2-r)}\left(e^{-r \ln r+i r \pi}-e^{-r \ln r-i r \pi}\right) d r \\
= & \frac{1}{-\pi x} \int_{0}^{\infty} \frac{(-x)^{r}}{(1-r) \Gamma(1-r)} \frac{\sin (r \pi)}{r^{r}} d r \\
= & \frac{1}{\pi(-x)} \int_{0}^{\infty} \frac{(-x)^{r}}{1-r} \frac{\Gamma(r) \sin (r \pi)}{\pi} \frac{\sin (r \pi)}{r^{r}} d r \\
= & \frac{1}{\pi^{2}(-x)} \int_{0}^{\infty} \frac{(-x)^{r}}{r^{r}} \sin ^{2}(r \pi) \frac{\Gamma(r)}{1-r} d r .
\end{aligned}
$$

Finally, we investigate the rise of $f$ and its asymptotic behavior in the points $-e$ and 0 .

Theorem 2.4. Let $x$ be in $(-e, 0)$; then:
(i) $f(x) \sim\left(\sqrt{2} /\left(e^{3 / 2} \pi\right)\right)(e+x)^{1 / 2}$, as $x \rightarrow-e_{+}$,
(ii) $f(x) \sim \ln ^{-2}(-x)(-x)^{-1}$, as $x \rightarrow 0_{-}$,
(iii) $f$ is increasing strictly on $(-e, 0)$.

Proof. (i) By Theorem 2.3(ii) and with $u:=1-\ln (-x)$ we get:

$$
f(x)=\frac{1}{2 \pi i} \int_{c-i \infty}^{c+i \infty} \frac{(s-1)^{s-1}}{\Gamma(s+1)} e^{-s} e^{s u} d s
$$

Moreover, with Stirling's formula (see [7, p. 294]), as $s \rightarrow \infty$, it follows that

$$
\begin{aligned}
\frac{(s-1)^{s-1}}{\Gamma(s+1)} e^{-s} & =\frac{(s-1)^{s-1}}{s(s-1) \Gamma(s-1)} e^{-s} \\
& \sim \frac{(s-1)^{s-1} e^{-s}}{s(s-1)(s-1)^{s-1} e^{-(s-1)}} \sqrt{\frac{s-1}{2 \pi}} \sim \frac{s^{-3 / 2}}{e \sqrt{2 \pi}} .
\end{aligned}
$$

Thus the conditions of an Abel theorem for inversions of Laplace transforms (Theorem 3 in [4, p. 503]) are fulfilled and as $u \rightarrow 0_{+}$and $x \rightarrow-e_{+}$ respectively $f$ holds:

$$
\begin{aligned}
f(x) & \sim \frac{1}{e \sqrt{2 \pi}} \frac{u^{1 / 2}}{\Gamma\left(\frac{3}{2}\right)}=\frac{\sqrt{2}}{e \pi}(1-\ln (-x))^{1 / 2} \\
& =\frac{\sqrt{2}}{e \pi}\left(-\ln \left(1-\frac{x+e}{e}\right)\right)^{1 / 2} \\
& \sim \frac{\sqrt{2}}{e \pi}\left(\frac{x+e}{e}\right)^{1 / 2}=\frac{\sqrt{2}}{e^{3 / 2} \pi}(x+e)^{1 / 2} .
\end{aligned}
$$

(ii) Using Theorem 2.3(iii), we have $f(x)=\tilde{f}(-\ln (-x)) /-x$, with

$$
\tilde{f}(-\ln (-x)):=\int_{0}^{\infty} e^{-(-\ln (-x)) r} \frac{\sin ^{2}(\pi r)}{\pi^{2} r^{r}} \frac{\Gamma(r)}{1-r} d r
$$

Since

$$
\frac{\sin ^{2}(\pi r)}{\pi^{2} r^{r}} \frac{\Gamma(r)}{1-r}=\frac{\sin (\pi r)}{\pi r^{r}} \frac{1}{\Gamma(2-r)} \sim r, \quad \text { as } \quad r \rightarrow 0_{+},
$$

we can apply an Abel theorem for Laplace integrals (Theorem 33.3 in [3, p. 241]) and obtain the asymptotic behavior of $\tilde{f}$ as $-\ln (-x) \rightarrow \infty$,

$$
\tilde{f}(-\ln (-x)) \sim \frac{\Gamma(1+1)}{(-\ln (-x))^{2}}, \quad \text { as } \quad-\ln (-x) \rightarrow \infty
$$

and that means $f(x) \sim \ln ^{-2}(-x)(-x)^{-1}$, as $x \rightarrow 0_{-}$.
(iii) Using Theorem 2.3 (iii) and a theorem on the derivative of Laplace integrals (Theorem 6.1 in [3, p.37]), it follows that

$$
f^{\prime}(x)=\frac{1}{(-x)^{2} \pi^{2}} \int_{0}^{\infty} \frac{(-x)^{r}}{r^{r}} \sin ^{2}(\pi r) \Gamma(r) d r
$$

which is greater than 0 for all $x$ in $(-e, 0)$, and thus the proof is completed.

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[^0]:    ${ }^{1}$ The result in (iii) is not surprising; compare with the sin-term in (1.2).

